

# Chapter 32 Inductance

## 32.1 Self-Induction and Inductance

### Self-Inductance

$$\Phi_m = BA = \mu_0 nIA \propto I \quad \rightarrow \quad \Phi_m = LI$$

The unit of the inductance is henry (H).  $1\text{H} = 1 \frac{\text{Wb}}{\text{A}} = 1 \frac{\text{T} \cdot \text{m}^2}{\text{A}}$

When the current in the circuit is changing, the magnetic flux is also changing.

$$\frac{d\Phi}{dt} = \frac{d(LI)}{dt} = L \frac{dI}{dt} \quad \rightarrow \quad \text{The induced emf should be } \varepsilon = -\frac{d\Phi}{dt} = -L \frac{dI}{dt}$$

The self inductance of a infinite long solenoid:

$$\Phi_m = NAB = NA\mu_0 \frac{N}{l} I = \mu_0 A \frac{N^2}{l} I \quad \rightarrow \quad L = \frac{\Phi_m}{I} = \mu_0 A \frac{N^2}{l}$$

$$\varepsilon = -L \frac{dI}{dt}$$

Considering the inductor having an internal resistance  $r$ , the potential difference is:

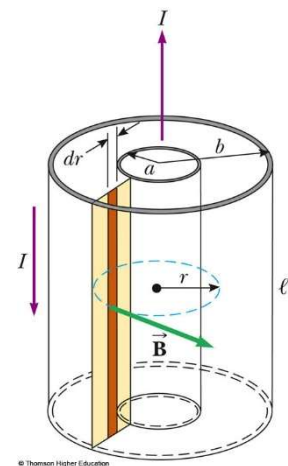
$$\Delta V = \varepsilon - Ir = -L \frac{dI}{dt} - Ir$$

Example: Model a long coaxial cable as two thin, concentric, cylindrical conducting shells of radii  $a$  and  $b$  and length  $l$ . The conducting shells carry the same current in opposite directions. Calculate the inductance  $L$  of this cable.

Calculate magnetic flux:

$$B = \frac{\mu_0 I}{2\pi r} \quad \rightarrow \quad \Phi = \int_0^l \int_a^b \frac{\mu_0 I}{2\pi r} dr dz = \frac{\mu_0 I}{2\pi} l \ln\left(\frac{b}{a}\right)$$

$$\Phi = LI \quad \rightarrow \quad L = \frac{\mu_0 l}{2\pi} \ln\left(\frac{b}{a}\right)$$



## 32.2 RL Circuits

Use Kirchhoff's rule:

$$\oint \vec{E} \cdot d\vec{l} = 0$$

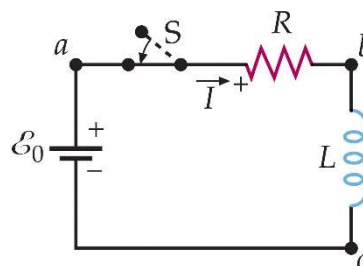
$$\varepsilon_0 - IR - L \frac{dI}{dt} = 0$$

$$\text{Differential Eq: } \varepsilon_0 - IR - L \frac{dI}{dt} = 0 \rightarrow \varepsilon_0 - IR = L \frac{dI}{dt}$$

$$\rightarrow dt = L \frac{dI}{(\varepsilon_0 - IR)} \rightarrow dt = -\frac{L}{R} \frac{d(\varepsilon_0 - IR)}{(\varepsilon_0 - IR)}$$

$$\rightarrow -\frac{R}{L}t = \ln\left(\frac{\varepsilon_0 - IR}{\varepsilon_0}\right) \rightarrow \varepsilon_0 - IR = \varepsilon_0 e^{-\frac{R}{L}t} \rightarrow I = \frac{\varepsilon_0}{R} \left(1 - e^{-\frac{R}{L}t}\right)$$

$$\text{Time constant: } \tau = \frac{L}{R}$$



Example: Find the total energy dissipated in the resistor R, when the current in the inductor decreases from its initial value of  $I_0$  to 0?

$$I = \frac{\varepsilon_0}{R} e^{-\frac{R}{L}t}, \quad P = I^2 R \rightarrow U = \int_0^{\infty} I_0^2 e^{-2\frac{R}{L}t} R dt = \frac{1}{2} LI_0^2$$

## 32.3 Energy in a Magnetic Field

Obtain the magnetic energy from the emf induced by self inductance.

$$\Phi_m = LI \rightarrow \text{The induced emf is } \varepsilon = -\frac{d\Phi_m}{dt} = -L \frac{dI}{dt}$$

$$\text{The energy dissipated or the power is } P = IV = I\varepsilon = -LI \frac{dI}{dt}$$

The total energy when the current has reached its final value  $I_f$  is:

$$U = \int_{t=0}^{t=I_f} \left( LI \frac{dI}{dt} \right) dt = \int_{I=0}^{I=I_f} LI dI = \frac{1}{2} LI^2$$

Calculate the magnetic energy by obtaining the energy stored in the self inductor of an infinite solenoid.

$$B = \mu_0 n I, \quad \Phi = n l A (\mu_0 n I) = L I$$

$$U = \frac{1}{2} L I^2 = \frac{1}{2} \mu_0 n^2 (A l) I^2 = u_m V$$

$$u_m = \frac{1}{2} \mu_0 n^2 I^2 = \frac{B^2}{2\mu_0} \quad \langle \text{----} \rangle \quad u_e = \frac{1}{2} \epsilon_0 E^2 \quad (\text{Do you remember how to get this?})$$

Example: A certain region of space contains a uniform magnetic field of 0.020 T and a uniform electric field of  $2.5 \times 10^6$  N/C. Find (a) the total electromagnetic density.

$$u_e = \frac{1}{2} \epsilon_0 E^2 = \frac{1}{2} (8.85 \times 10^{-12}) (2.5 \times 10^6)^2 = 27.7 \text{ J/m}^3$$

$$u_m = \frac{B^2}{2\mu_0} = \frac{(0.02)^2}{2(4\pi \times 10^{-7})} = 159 \text{ J/m}^3$$

## 32.4 Mutual Inductance

### Mutual Inductance

The magnetic field of loop 1 is:  $B \sim \frac{\mu_0 I_1}{2R}$

The flux at 2 is  $\Phi_2 \sim \frac{\mu_0 I_1}{2R} \pi R^2 = \mu_0 \frac{\pi R^2}{2} I_1 = I_1 M_{12}$



The flux at 1 is  $\Phi_1 \sim \mu_0 \frac{\pi R^2}{2} I_2 = I_2 M_{21}$

The concept of inductance:  $\Phi = M I$

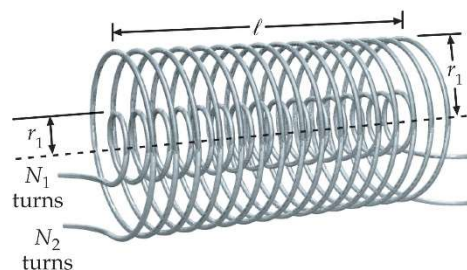
$M_{12} = M_{21}$  --> The mutual inductance is determined when the geometrical configuration between the two loops is given.

B in 1 due to 2:  $B = \mu_0 \frac{N_2}{l} I_2$

Flux in 1:  $\Phi = N_1 (\pi r_1^2) \left( \mu_0 \frac{N_2}{l} I_2 \right) = I_2 M_{21}$

B in 2 due to 1:  $B = \mu_0 \frac{N_1}{l} I_1$

Flux in 2:  $\Phi = N_2 (\pi r_2^2) \left( \mu_0 \frac{N_1}{l} I_1 \right) = I_1 M_{12}$



$$M_{12} = M_{21} = \mu_0 (\pi r_1^2) \frac{N_1 N_2}{l}$$

## 32.5 Oscillations in an LC Circuit

Kirchhoff's Rule:

$$-L \frac{dI}{dt} - \frac{Q}{C} = 0 \rightarrow L \frac{d^2 Q}{dt^2} + \frac{Q}{C} = 0$$

Compare with:  $F = ma = -kx$

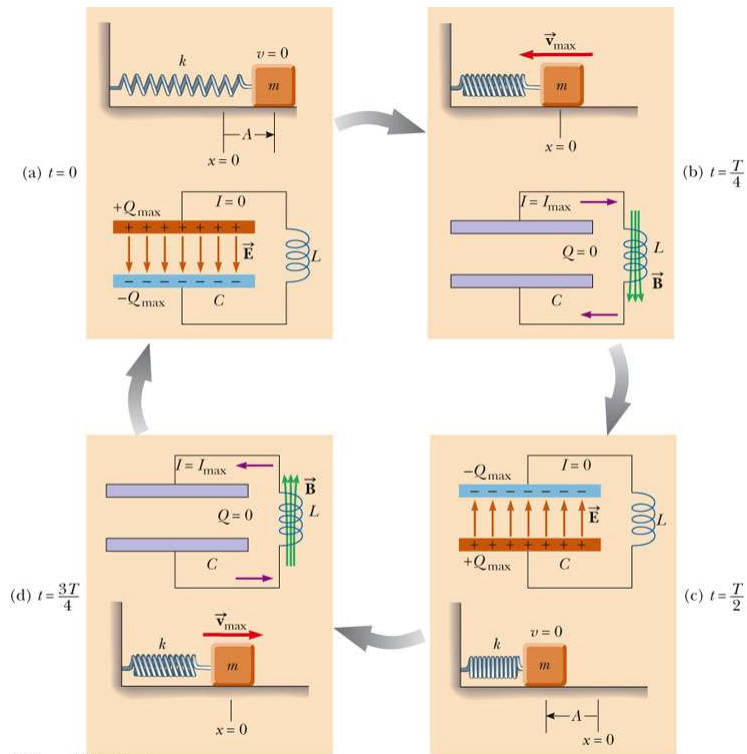
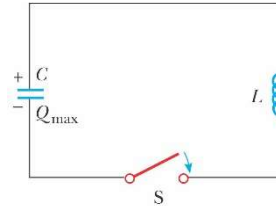
$$m \frac{d^2 x}{dt^2} + kx = 0 \quad \text{solution: } x(t) = A \cos(\omega t + \phi)$$

$$L \frac{d^2 Q}{dt^2} + \frac{Q}{C} = 0 \quad \text{solution: } Q(t) = Q_{\max} \cos(\omega t + \phi), \quad \omega = \frac{1}{\sqrt{LC}}$$

$$\rightarrow I = -\omega Q_{\max} \sin(\omega t + \phi)$$

$$U = U_C + U_L = \frac{Q^2}{2C} + \frac{L}{2} I^2 = \frac{1}{2C} Q_{\max}^2 \cos^2(\omega t + \phi) + \frac{L}{2} \omega^2 Q_{\max}^2 \sin^2(\omega t + \phi)$$

$$= \frac{Q_{\max}^2}{2C}$$

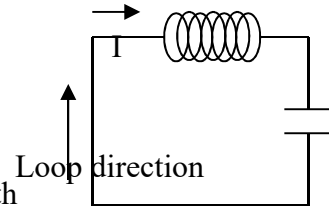


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Apply the Kirchoff's loop rule:

$$-L \frac{dI}{dt} - \frac{Q}{C} = 0$$

$$\rightarrow L \frac{d^2Q}{dt^2} + \frac{Q}{C} = 0 \quad (2^{\text{nd}} \text{ Differential Equation, compare with}$$



harmonic oscillation:  $F = ma = m \frac{d^2x}{dt^2} = -kx$  with the answer of  $x = A \cos(\omega t + \delta)$ ,

where  $\omega = \sqrt{\frac{k}{m}}$

Remember the pattern of this differential equation:

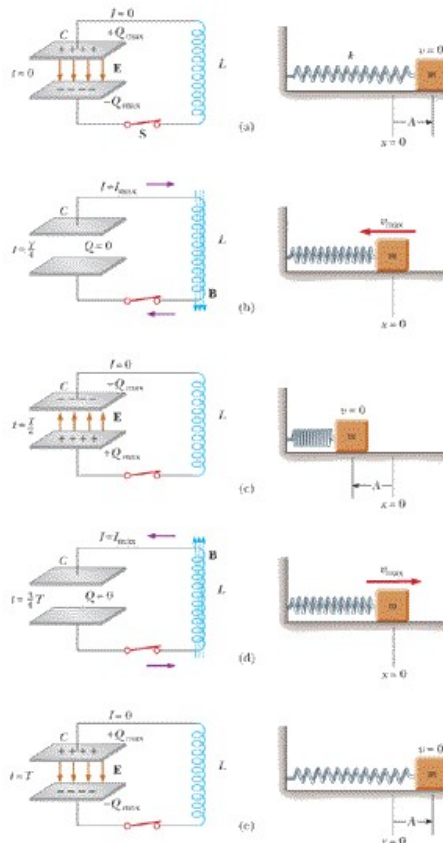
$L \frac{d^2Q}{dt^2} + \frac{Q}{C} = 0$  --> the solutions are periodical functions and the useable functions are  $\sin(x)$  ( $\cos(x)$ ),  $\exp(ix)$  ( $\exp(x)$ ), ...

Guess that the answer is  $Q = A \cos(Bt + C)$ . (Here  $A$  and  $C$  can be determined by initial conditions.)

$$\rightarrow -LB^2 A \cos(Bt + C) + \frac{1}{C} A \cos(Bt + C) = 0 \rightarrow B = \frac{1}{\sqrt{LC}} \equiv \omega$$

$$\rightarrow Q = A \cos(\omega t - \delta) \quad \& \quad I = \frac{dQ}{dt} = -\omega A \sin(\omega t - \delta) = -I_{\text{peak}} \sin(\omega t)$$

Serway/Jewett; Principles of Physics, 3/e  
Figure 24.9



Capacitor --> Electric Field  
--> Potential Energy

Inductor --> Moving of  
Charges --> Kinetic Energy

The average energy stored in the capacitor (inductor) is  $\frac{Q^2}{2C}$  ( $\frac{1}{2}LI^2$ ).

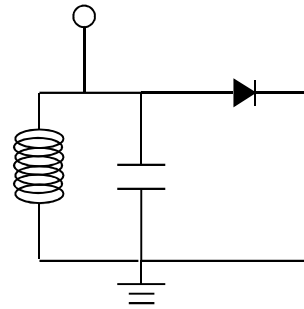
The instantaneous energy transferring in the circuit is:

$$\frac{Q^2}{2C} + \frac{1}{2}LI^2 = \frac{A^2 \cos^2(\omega t - \delta)}{2C} + \frac{1}{2}L\omega^2 A^2 \sin^2(\omega t - \delta) = \frac{A^2}{2C} = \frac{Q_{peak}^2}{2C} = \frac{1}{2}LI_{peak}^2$$

What are physical pictures of  $Q_{peak}$  and  $I_{peak}$  ?

Example: A 2- $\mu$ F capacitor is charged to 20 V and the capacitor is then connected across a 6- $\mu$ H inductor. (a) What is the frequency of oscillation? (b) What is the peak value of the current?

$$(a) \omega = \frac{1}{\sqrt{LC}}, (b) \frac{CV^2}{2} = \frac{Q_{peak}^2}{2C} = \frac{LI_{peak}^2}{2}$$



Simple AM Radio receiver:

## 32.6 The RLC Circuit

Kirchhoff's Rule:

$$-IR - L \frac{dI}{dt} - \frac{Q}{C} = 0$$

Power Consideration:

$$P = LI \frac{dI}{dt} + \frac{Q}{C} I = -I^2 R$$

$$\rightarrow L \frac{d^2 Q}{dt^2} + R \frac{dQ}{dt} + \frac{Q}{C} = 0$$

Compare with damped oscillation:

$$m \frac{d^2 x}{dt^2} + b \frac{dx}{dt} + kx = 0 \rightarrow x(t) = Ae^{nt}, mn^2 + bn + k = 0 \rightarrow n = \frac{-b \pm \sqrt{b^2 - 4mk}}{2m}$$

$$b^2 < 4mk \rightarrow x(t) = Ae^{-\frac{b}{2m}t} \cos\left(\frac{\sqrt{4mk - b^2}}{2m}t\right)$$

$$Q(t) = Q_{max} e^{-\frac{R}{2L}t} \cos\left(\frac{\sqrt{4L/C - R^2}}{2L}t\right)$$

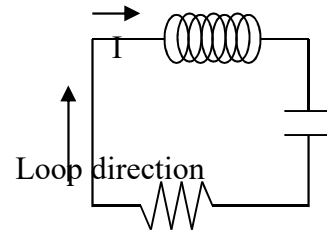
**TABLE 32.1**

**Analogies Between Electrical and Mechanical Systems**

| Electric Circuit   |  | One-Dimensional Mechanical System                 |
|--|--|---|
| Charge   | $Q \leftrightarrow x$  | Position  |
| Current  | $I \leftrightarrow v_x$  | Velocity  |
| Potential difference   | $\Delta V \leftrightarrow F_x$   | Force   |
| Resistance   | $R \leftrightarrow b$  | Viscous damping coefficient                       |
| Capacitance  | $C \leftrightarrow 1/k$  | ( $k$ = spring constant)                          |
| Inductance   | $L \leftrightarrow m$  | Mass  |
| Current = time derivative of charge                          | $I = \frac{dQ}{dt} \leftrightarrow v_x = \frac{dx}{dt}$  | Velocity = time derivative of position            |
| Rate of change of current = second time derivative of charge | $\frac{dI}{dt} = \frac{d^2Q}{dt^2} \leftrightarrow a_x = \frac{dv_x}{dt} = \frac{d^2x}{dt^2}$                        | Acceleration = second time derivative of position |
| Energy in inductor   | $U_L = \frac{1}{2}LI^2 \leftrightarrow K = \frac{1}{2}mv^2$  | Kinetic energy of moving object                   |
| Energy in capacitor  | $U_C = \frac{1}{2}\frac{Q^2}{C} \leftrightarrow U = \frac{1}{2}kx^2$   | Potential energy stored in a spring               |
| Rate of energy loss due to resistance                        | $I^2R \leftrightarrow bv^2$  | Rate of energy loss due to friction               |
| RLC circuit  | $L\frac{d^2Q}{dt^2} + R\frac{dQ}{dt} + \frac{Q}{C} = 0 \leftrightarrow m\frac{d^2x}{dt^2} + b\frac{dx}{dt} + kx = 0$ | Damped object on a spring                         |

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## RLC Circuit (Damped Oscillation)



$$\text{Diff Eq: } -L \frac{dI}{dt} - \frac{Q}{C} - IR = 0$$

$$\rightarrow L \frac{d^2Q}{dt^2} + R \frac{dQ}{dt} + \frac{Q}{C} = 0 \quad (\text{Damped Oscillations: } F = ma = -kx - bv)$$

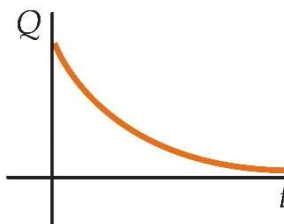
The natural frequency (no resistor) is:  $\omega_0 = \frac{1}{\sqrt{LC}}$  ( $L \frac{d^2Q}{dt^2} + \frac{Q}{C} = 0$ ).

We guess a solution of  $Q = Ae^{Bt}$  for solving the differential equation

$$L \frac{d^2Q}{dt^2} + R \frac{dQ}{dt} + \frac{Q}{C} = 0.$$

$$\left( LB^2 + RB + \frac{1}{C} \right) Ae^{Bt} = 0 \rightarrow B = \frac{-R \pm \sqrt{R^2 - 4\frac{L}{C}}}{2L} = -\frac{R}{2L} \pm \sqrt{\left(\frac{R}{2L}\right)^2 - \frac{1}{LC}}$$

over-damped:  $\left(\frac{R}{2L}\right)^2 - \frac{1}{LC} > 0$



under-damped:  $\left(\frac{R}{2L}\right)^2 - \frac{1}{LC} < 0$

under-damped solution:  $Q = Ae^{-\frac{R}{2L}t} e^{\pm i\sqrt{\omega_0^2 - \left(\frac{R}{2L}\right)^2}t}$

The energy distributed in the circuit elements is:

$$L \frac{d^2Q}{dt^2} + R \frac{dQ}{dt} + \frac{Q}{C} = 0 \quad \rightarrow \quad LI \frac{dI}{dt} + RI^2 + \frac{Q}{C} \frac{dQ}{dt} = 0$$

$$\rightarrow \frac{d}{dt} \left( \frac{1}{2} LI^2 + \frac{Q^2}{2C} \right) + I^2 R = 0$$

